Phase Space Factors

For quantum transitions to un-bound states — for example, an atom emitting a photon, or a radioactive decay, or scattering — which is a kind of unbound → unbound transition, — the transition rate is given by the Fermi’s golden rule:

\[
\Gamma \overset{\text{def}}{=} \frac{d \text{probability}}{d \text{time}} = \frac{2\pi \rho}{\hbar} \times \left| \langle \text{final} | \hat{T} | \text{initial} \rangle \right|^2
\]  

(1)

where \( \hat{T} = \hat{H}_\text{perturbation} + \text{higher order corrections} \), and \( \rho \) is the density of final states,

\[
\rho = \frac{dN_{\text{final states}}}{dE_{\text{final}}}.
\]  

(2)

For an example, consider an atom emitting a photon of specific polarization \( \lambda \). Using the large-box normalization for the photon’s states, we have

\[
dN_{\text{final}} = \left( \frac{L}{2\pi} \right)^3 d^3 k_\gamma = \frac{L^3}{(2\pi)^3} \times k_\gamma^2 dk_\gamma d^2 \Omega_\gamma
\]  

(3)

while \( dE_{\text{final}} = dE_\gamma = \hbar c \times dk_\gamma \), hence

\[
\rho = L^3 \times \frac{k_\gamma^2}{(2\pi)^3 \hbar c} \times d^2 \Omega_\gamma.
\]  

(4)

The \( L^3 \) factor here cancels against the \( L^{-3/2} \) factor in the matrix element \( \langle \text{atom'} + \gamma | \hat{T} | \text{atom} \rangle \) due to the photon’s wave function in the large-box normalization. As to the remaining \( d^2 \Omega_\gamma \) factor, we should integrate over it to get the total decay rate, or divide by it to get the partial emission rate \( d\Gamma/d\Omega \) for the photons emitted in a particular direction, thus

\[
\frac{d\Gamma(\lambda)}{d\Omega} = \frac{k^2}{(2\pi)^3 \hbar c} \times L^3 \left| \langle \text{atom'} + \gamma(k, \lambda) | \hat{T} | \text{atom} \rangle \right|^2,
\]

\[
\Gamma_{\text{total}} = \int d\Omega \sum_\lambda \frac{k^2}{(2\pi)^3 \hbar c} \times L^3 \left| \langle \text{atom'} + \gamma(k, \lambda) | \hat{T} | \text{atom} \rangle \right|^2.
\]  

(5)

In relativistic normalization of quantum states and matrix elements, there are no \( L^{-3/2} \) factors but instead there are \( \sqrt{2E} \) factors for each final-state or initial state particle, and
they must be compensated by dividing the density of states \( \rho \) by the \( \prod E_i \). Also, we must allow for motion of all the final-state particles (i.e., both the photon and the recoiled atom) but impose the momentum conservation as a constraint. Thus, for a decay of 1 initial particle into \( n \) final particles,

\[
\Gamma = \frac{1}{2E_i} \int \frac{d^3p'_1}{(2\pi)^3 2E'_1} \cdots \int \frac{d^3p'_n}{(2\pi)^3 2E'_n} |\langle p'_1, \ldots, p'_n | M | p_{in} \rangle|^2 \times (2\pi^4) \delta(4)(p'_1 + \cdots + p'_n - p_{in}),
\]

(6)

where the \( \delta \) function takes care of both momentum conservation and of the denominator \( dE_f \) in the density-of-states factor (2). Likewise, the transition rate for a generic \( 2 \rightarrow n \) scattering process is given by

\[
\Gamma = \frac{1}{2E_1 \times 2E_2} \int \frac{d^3p'_1}{(2\pi)^3 2E'_1} \cdots \int \frac{d^3p'_n}{(2\pi)^3 2E'_n} |\langle p'_1, \ldots, p'_n | M | p_{1,2} \rangle|^2 \times (2\pi^4) \delta(4)(p'_1 + \cdots + p'_n - p_1 - p_2).
\]

(7)

In terms of the scattering cross-section \( \sigma \), the rate (7) is \( \Gamma = \sigma \times \text{flux of initial particles} \). In the large-box normalization the flux is \( L^{-3} |v_1 - v_2| \), so in the continuum normalization it’s simply the relative speed \( |v_1 - v_2| \). Consequently, the total scattering cross-section is given by

\[
\sigma_{tot} = \frac{1}{4E_1 E_2 |v_1 - v_2|} \int \frac{d^3p'_1}{(2\pi)^3 2E'_1} \cdots \int \frac{d^3p'_n}{(2\pi)^3 2E'_n} |\langle p'_1, \ldots, p'_n | M | p_{1,2} \rangle|^2 \times (2\pi^4) \delta(4)(p'_1 + \cdots + p'_n - p_1 - p_2).
\]

(8)

In particle physics, all the factors in eqs (6) or (8) besides the matrix elements — as well as the integrals over such factors — are collectively called the phase space factors.

A note on Lorentz invariance of decay rates or cross-sections. The matrix elements \( \langle \text{final} | M | \text{initial} \rangle \) are Lorentz invariant, and so are all the integrals over the final-particles’ momenta and the \( \delta \)-functions. The only non-invariant factor in the decay-rate formula (6) is the pre-integral \( 1/E_{init} \), hence the decay rate of a moving particle is

\[
\Gamma(\text{moving}) = \Gamma(\text{rest frame}) \times \frac{M}{E}
\]

(9)

where \( M/E \) is precisely the time dilation factor in the moving frame.
As to the scattering cross-section, it should be invariant under Lorentz boosts along the initial axis of scattering, thus the same cross-section in any frame where \( \mathbf{p}_1 \parallel \mathbf{p}_2 \). This includes the lab frame where one of the two particles is initially at rest, the center-of-mass frame where \( \mathbf{p}_1 + \mathbf{p}_2 = 0 \), and any other frame where the two particles collide head-on. And indeed, the pre-integral factor in eq. \((8)\) for the cross-section

\[
\frac{1}{4E_1 E_2 |\mathbf{v}_1 - \mathbf{v}_2|} = \frac{1}{4|\mathbf{E}_2 \mathbf{p}_1 - \mathbf{E}_1 \mathbf{p}_2|}
\]

is invariant under Lorentz boosts along the scattering axis.

Let’s simplify eq. \((8)\) for a 2-particle \(\rightarrow\) 2-particle scattering process in the center-of-mass frame where \(\mathbf{p}_1 + \mathbf{p}_2 = 0\). In this frame, the pre-exponential factor \((10)\) becomes

\[
\frac{1}{4 \mathbf{p} \times (E_1 + E_2)}
\]

while the remaining phase space factors amount to

\[
P_{\text{int}} = \int \frac{d^3 \mathbf{p}'}{(2\pi)^3} \int \frac{d^3 \mathbf{p}'}{(2\pi)^3} (2\pi)^4 \delta(3)(\mathbf{p}_1' + \mathbf{p}_2') \delta(E_1' + E_2' - E_{\text{net}})
\]

\[
= \int \frac{d^3 \mathbf{p}'}{(2\pi)^3} \times 2E_1' \times 2E_2' \delta(E_1'(\mathbf{p}_1') + E_2'(-\mathbf{p}_1') - E_{\text{net}})
\]

\[
= \int d^2 \Omega_{\mathbf{p}'} \times \int dp' \frac{p'^2}{16\pi^2 E_1'E_2'} \times \delta(E_1' + E_2' - E_{\text{tot}})
\]

\[
= \int d^2 \Omega_{\mathbf{p}'} \left[ \frac{p'^2}{16\pi^2 E_1'E_2'} \right] \text{ when } E_1' + E_2' = E_{\text{tot}}.
\]

On the last 3 lines here \(E_1' = E_1'(\mathbf{p}_1') = \sqrt{p'^2 + m_1^2}\) while \(E_2' = E_2'(\mathbf{p}_2') = -\mathbf{p}_1' = \sqrt{p'^2 + m_2^2}\. Consequently,

\[
\frac{dE_1'}{dp'} = \frac{p'}{E_1'}, \quad \frac{dE_2'}{dp'} = \frac{p'}{E_2'},
\]

hence

\[
\frac{d(E_1' + E_2')}{dp'} = \frac{p'}{E_1'} + \frac{p'}{E_2'} = \frac{p'}{E_1'E_2'} \times (E_2' + E_1' = E_{\text{tot}}),
\]

\[3\]
and therefore
\[ \mathcal{P}_{\text{int}} = \frac{1}{16\pi^2} \times \frac{p'}{E_{\text{tot}}} \times \int d^2\Omega_{p'}. \]  
(15)

Including the pre-integral factor (11), we arrive at the net phase space factor
\[ \mathcal{P} = \frac{p'}{p} \times \frac{1}{64\pi^2 E_{\text{tot}}^2} \times \int d^2\Omega_{p'}. \]  
(16)

The matrix element \( \mathcal{M} \) for the scattering should be put inside the direction-angle integral in this phase-space formula. Thus, the total scattering cross-section is
\[ \sigma_{\text{tot}}(1 \rightarrow 1' + 2') = \frac{p'}{p} \times \frac{1}{64\pi^2 E_{\text{cm}}^2} \times \int d^2\Omega \ |\langle p' + p_2 | \mathcal{M} | p_1 + p_2 \rangle|^2, \]  
(17)

while the partial cross-section for scattering in a particular direction is
\[ \frac{d\sigma(1 \rightarrow 1' + 2')}{d\Omega_{\text{cm}}} = \frac{p'}{p} \times \frac{1}{64\pi^2 E_{\text{cm}}^2} \times |\langle p' + p_2 | \mathcal{M} | p_1 + p_2 \rangle|^2. \]  
(18)

Note: the total cross-section is the same in frames where the initial momenta are collinear, but in the partial cross-section, \( d\Omega \) depends on the frame of reference, so eq. (18) applies only in the center-of-mass frame. Also, the \( E_{\text{cm}} \) factor in denominators of both formulae stands for the net energy in the center-of-mass frame. In frame-independent terms,
\[ E_{\text{cm}}^2 = (p_1 + p_2)^2 = (p_1' + p_2')^2 = s. \]  
(19)

Finally, let me write down the phase-space factor for a 2-body decay (1 particle \( \rightarrow \) 2 particles) in the rest frame of the initial particle. The under-the-integral factors for such a decay are the same as in eq. (15) for a \( 2 \rightarrow 2 \) scattering, but the pre-integral factor is \( 1/2M_{\text{init}} \) instead of the (11), thus
\[ \mathcal{P} = \frac{p'}{32\pi^2 M^2}, \]  
(20)

meaning
\[ \frac{d\Gamma(0 \rightarrow 1' + 2')}{d\Omega} = \frac{p'}{32\pi^2 M^2} \times |\langle p' + p_2 | \mathcal{M} | p_0 \rangle|^2, \]  
(21)

\[ \Gamma(0 \rightarrow 1' + 2') = \frac{p'}{32\pi^2 M^2} \times \int d^2\Omega \ |\langle p' + p_2 | \mathcal{M} | p_0 \rangle|^2. \]  
(22)